The CMS Silicon Tracker:
 from the performance in cosmic runs to the p_T resolution in early
 data

Candidata: Maria Assunta Borgia Relatore: Prof. Marco Costa Controrelatore: Dott. Antonio Pellegrino



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The CMS Silicon Strip Tracker

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Tracker commissioning with cosmic muons

> Measurement of the p_T resolution from the Z line-shape

Introduction on the method

Ideal scenario

Early data scenario

Data driven method for the correction of the MC resolution

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Why LHC?



- Energy scales never explored until now... new physics foreseen!
- Higgs search and Electroweak symmetry breaking: crucial points for Standard Model
- Early data: re-discovery of the Standard Model physics (e.g. $Z \rightarrow \mu\mu$, $W \rightarrow \mu\nu$) as
 - Background to new physics searches
 - Fundamental channels to improve PDF knowledge
 - Detector calibration: resolution measurement and momentum scale correction
 - Tuning of Monte Carlo simulation

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The Compact Muon Solenoid

- Compact cylindrical detector, "onion" structure:
 - Muon Chambers (muons trigger)
 - Hadronic + Electromagnetic calorimeters-
 - Silicon Strip Tracker -
 - Pixel Vertex Detector
- Very strong magnetic field (4T) to bend the particles trajectory and permit a precise measurement of the transverse momentum







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The CMS Tracker: layout

Physics Enviroment	Design Requirements
High particle fluence	Radiation hardness
High track density	High granularity
25 ns bunch crossing	Fast read-out



- Tracking system designed to provide a precise and efficient measurement of the charged particle trajectories in the LHC collisions
 - $\mathbf{B} = \mathbf{4}$ Tesla
 - **Resolution:** $\Delta p_t/p_t \sim 1-2\%$ ($|\eta| < 1.6$)
 - Tracking efficiency: ε~99% (μ), ~90% hadrons
- Silicon Pixels surrounded by Silicon Strip detectors
 - Pixels:
 - $\sim 1 \text{ m}^2$ of Si sensors, 65 M channels, 100x150 mm², r = 4, 7, 11 cm
 - Strips:
 - ~ 200 m² of Si sensors, 15148 modules, ~10 M channels
 - 10 barrel layers (TIB,TOB)
 - 12 end-cap wheels per side (TID,TEC)

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The CMS Tracker: working principle



The Tracker calibration workflow

Before starting to collect physics data, the detector needs to be fully commissioned and calibrated in order to have the most reliable measurements to be used for physics studies.

- **Gain calibration**
- Noise measurement
- **Cluster properties studies**
- **Identification of bad components**
- Lorentz angle calibration
- **Tracking performance**
- Alignment of silicon modules R. Castello's presentation

Data Quality Monitoring system

- The Data Quality Monitoring (DQM) constantly controls the status of the detector and of the reconstruction
 - capable of running on a variety of online and offline environments, in the control room as well as in remote sites
 - Monitored quantities
 - Read-Out (Fed errors), Raw Data, Cluster/Hit Properties, Tracks
 - Calibration constants used during the reconstruction
 - Quality test applied to quickly spot problems during online/offline reconstruction
 - Fundamental tool for the calibration



The Tracker Slice Test @ TIF CRuZeT/CRAFT @ P5

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The Tracker Slice Test

First large scale system test: Tracker Integration Facility (TIF) @ CERN (ground level).

- ~15% of the full detector operated from Feb. to July '07
- *a* five operating temperatures (15, 10, -1, -10, -15 °C)
- More than 4 M events
- Verified HW, SW and calibration procedures in conditions close to the final one



Trigger:

- rate: ~6 Hz
- Flexible trigger geometry
- 5 cm lead on bottom scintillator to reject soft muons

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CRuZeT & CRAFT

Tracker inserted in CMS in December 2007 and fully cabled and tested in the following months

In Global Run with the full CMS in June 2008

- 98% of the detector switched on
- Temperature of operation 10°
- ~8.5 M cosmic data with and without B field
- Magnetic field OFF in CRuZeT (Cosmic Run @ Zero Tesla) and 3.8 T in CRAFT (Cosmic Run @ Four Tesla)
- Cosmic Trigger configured using muon chambers

Gain calibration

- Aim: to achieve a uniform electronic gain among the readout chips (APV)
- Measurement of the value of the height of synchronization pulses, referred to as tick-marks, generated by the APV
- Equalization of the response of the full electronic chain to a known value (640 ADC)
- Influenced by voltage, temperature → important to monitor stability in time
 TIF results



Noise measurement

- Before irradiation, completely determined by the input capacitance load of the APV chips, dominated by the silicon strips.
- Linear dependence on the length of the silicon strips.
- When the modules are mounted on the final support structures (close proximity to each other), other possible sources of noise can arise (grounding loops, cross talk, digital noise, cables...), affecting the final noise performance.



Bad components identification

Online bad component identification:

- **•Search for HV errors**
- **•Search for noisy components**

Offline bad component identification:

Search for high occupancy strips or APV
 Search for inefficient components





Non gaussian behaving strips: HV off

- Known from commissioning
- Online identification
- **Offline** identification
- Online + offline
- **Currently: 2.7% of strips masked** 15

Cluster/hit properties studies

- Signal-to-noise renormalized to the detector thickness EXCELLENT parameter to monitor the stability of the tracker during data taking
- S_{ren} = S/K where K particle path length in the silicon. Landau distribution.
- $N = \sqrt{\sum_i N_i^2 / n_{strips}}$. Gaussian distribution.
- Independent on gain
- Sren/N largest possible
 - in order to reduce fake hits rate
 - will get worse for radiation damage





Fit: Landau 🕀

Gaussian

wrong synchronization with the trigger caused low S/N

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Lorentz angle calibration



- if **B** = **0** charges drift in **E** direction 0
- if $B \neq 0$ diffusion path influenced by B 0
- Net effect \Rightarrow
 - Shift in the hit position 0
 - Change in the cluster width Ο

$$< clusterwidth >= a + \left|\frac{t}{p} \cdot b \cdot (\tan \theta_t - \tan \Theta_L)\right|$$

where

• t = detector thickness

 \circ **p** = detector pitch

• a, b = coefficients expressing the carrier diffusion and the electronic capacitive coupling between nearby channels

$$\odot \Theta_{\rm L} = \text{Lorentz angle}$$

 $\circ \theta_t =$ track angle



3380

1 88

-0.1464

71.11/76

0.6

z

p'

n

track

Track properties (I): TIF

- **Three tracking algorithm, similar performances**
- **Number of tracks, number of hits per track,** χ^2 of the track, η , ϕ
- Monitoring stability in time



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Track properties (II): CRAFT

Magnetic field $ON \rightarrow$

• first possible estimate of the **p**_T of the particle!!!

- first possible estimate of the p_T resolution!!!
 - Tracks pointing to the interaction vertex
 - Track split in two legs
 - Projection of track parameters at the point of closest approach to the beam pipe
 - Residuals calculated: $\delta x = (x_1 x_2)/\sqrt{2}$
- Measurement of the tracking efficiency with inside-out method (as in LHC collision!)



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Measurement of p_T resolution from the Z line-shape

... and data driven correction of the Monte Carlo resolution on p_T

Measurement of p_T resolution from early data

Basic idea

Comparison between the observed invariant mass lineshape and the expected one:

 \circ Mass lineshape width \rightarrow measurement of the resolution

 \circ Mass peak shift \rightarrow bias in the momentum scale

Dependence on the track parameters

probabilistic approach necessary.

Estimation directly from data using a likelihood minimization

$$-\ln L = -\sum_{k=1}^{N_{ev}} \ln(P(m_k, s_k))$$

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Resolution measurement

Input:

 probability distributions f(M_{res},σ_{Mres})
 ansatz functions for resolution, scale (vs muon kinematics)



For each event the algorithm

otakes the reconstructed mass *m* with its resolution *s* estimated from the ansatz function

ocalculates the probability associated to these two values (m, s) using the probability distributions $f(M_{res}, \sigma_{Mres})$

ouse this probability value to construct the likelihood

ominimizes the likelihood with respect to the parameters of the ansatz functions.

Output: full set of parameters for resolution and scale

Samples and cuts

- An ideal MC $Z \rightarrow \mu\mu$ sample generated with Sherpa (CTEQ61) and reconstructed in an ideal scenario
- A "fake data" $Z \rightarrow \mu\mu$ sample from the same generated sample reconstructed in a realistic scenario after <u>10 pb⁻¹</u> with:
 - misalignment scenario after CRAFT alignment (the one at the startup)
 - real tracker noise condition, using noise values measured during CRAFT
 - list of bad component found during CRAFT excluded from the reconstruction
- Global muons (muons reconstructed in tracker + muon chambers) \rightarrow inner tracker track p_t measurement
- Only muons with $p_t > 20$ GeV (too poor statistics below)
- \sim 5000 Z survive after these cuts, after 10 pb⁻¹ data collected

Resolution measurement on ideal MC

- Momentum scale correction applied (no strong bias, linear in p_T)
- Resolution Function: basic function
 - constant in p_T
 - by point in η (single muon simulation)
 - 2 parameters



Comparison of the resolution from the fit of the ansatz function and the usual MC truth resolution (pt_{reco}-pt_{gen})/pt_{gen}. Excellent agreement between resolutions vs p_T and η from MC truth and from the algorithm



Resolution measurement on "fake data"

- Momentum scale correction more complicated:
 - linear in p_T
 - parabolic in η
 - o sinusoidal in φ
- Resolution Function:
 - \circ linear in p_T
 - \circ parabolic in η , separated regions for barrel and endcaps
 - 7 parameters

$$\sigma_{pT}/p_T = b_0 imes p_T + \left\{ egin{array}{ccc} b_1 imes \eta^2 & |\eta| > 0.6 \ \& \ |\eta| < 1.3 \ b_2 imes (|\eta| - b_3)^2 & \eta > 1.3 \ b_4 imes (|\eta| - b_5) + b_6 imes (|\eta| - b_5)^2 & \eta < -1.3 \end{array}
ight.$$





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pt(GeV)

Correction of the MC resolution

Basic idea: correct the MC resolution on muon momentum in order to reproduce the resolution measured on a real data sample \Rightarrow fundamental to compare data with a realistic MC!

Hypothesis 1: all the effects introduced in the reconstruction by a real condition detector wrt the ideal detector can be summarized in an additional gaussian smearing

Hypothesis 2: uncorrelated gaussians (\Rightarrow resolutions):

$$\sigma_{data}^{2} = \sigma_{MC}^{2} + \sigma_{add}^{2}$$

Smearing the muon transverse momentum by a Gaussian function:

$$p_t' = p_t * G(1, \sqrt{(\sigma_{data}^2 - \sigma_{MC}^2)})$$
 (1)

Resolution measurement on smeared MC



Resolution on IDEAL sample

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Resolution measurement on smeared MC



Resolution on REALISTIC sample

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Resolution measurement on smeared MC



The resolution ansatz function is capable to reproduce resolution from reco – gen comparison also for smeared MC events!

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Comparison between ideal MC, "fake data" and smeared MC: σ_{pT} vs p_t and η



Effect on the Z kinematics



Evident broadening of the Z di-muon mass in the case of fake data and smeared MC with respect to ideal MC!

No visible effect on Z transverse momentum

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Conclusions

- The tracker has been described in all its aspects: from the calibration and local reconstruction, to the effects that residual miscalibration and misalignment can have on the physics measurements.
- A big effort has been done in the last few years to ensure the most reliable physics measurements to be done in the very early phase of the data taking.
- Already with 10 pb⁻¹ data we are able to correct the muon momentum scale and to measure resolution on transverse momentum.
- The smearing method seems to reproduce quite well the "data" behavior starting from ideal MC ⇒ fundamental in order to have a realistic Monte Carlo!!
- Looking forward for the high energy collisions!!!

Thank you for your attention!



Likelihood construction

Decay muons cross detector material ⇒ measured muon kinematics bring to a reconstructed Z mass which is a resolution-smeared view of the B-W line-shape

$$P(m,s) = \int dx \frac{\sigma(x)}{s} e^{-\frac{(x-m)^2}{2s^2}} \quad (\sigma(\mathbf{x}) = \mathbf{Z} \text{ mass B-W})$$

This probability is used to construct a likelihood function:

$$-\ln L = -\sum_{k=1}^{N_{ev}} \ln(P_s(m_k, s_k) + P_b(m_k))$$

signal bkg

where *m* and *s* are parameterized in function of the observable quantities**:

$$m = m(P_{T,1}^{corr}, \phi_1, \cot \theta_1; P_{T,2}^{corr}, \phi_2, \cot \theta_2)$$
 measured mass

$$s = \sqrt{(\frac{\partial m}{\partial P_T})^2 \sigma_{P_T}^2 + (\frac{\partial m}{\partial \phi})^2 \sigma_{\phi}^2 + (\frac{\partial m}{\partial \cot \theta})^2 \sigma_{\cot \theta}^2}$$
 expected resolution

**
$$E = \sqrt{P_T^2(1 + \cot^2 \theta) + m_\mu^2}; P_x = P_T \cos \phi; P_y = P_T \sin \phi; P_z = P_T \cot \theta.$$

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Scale correction and resolution measurement

Hypothesis on a scale factor:

$$P_T^{corr} = F(\vec{x}; \vec{\alpha}) P_T$$

- $\blacksquare \quad F \rightarrow ansatz \ function \ for \ the \ momentum \ scale$
- $x \rightarrow$ kinematic variables on which the momentum scale depends
- $\alpha \rightarrow$ parameters of the scale correction function, to be extracted from the likelihood minimization

$$s = \sqrt{(\frac{\partial m}{\partial P_T})^2 \sigma_{P_T}^2 + (\frac{\partial m}{\partial \phi})^2 \sigma_{\phi}^2 + (\frac{\partial m}{\partial \cot \theta})^2 \sigma_{\cot \theta}^2} \quad \text{where} \quad \begin{array}{l} \sigma_{P_T} = G_1(\vec{x};\vec{\beta}); \\ \sigma_{\phi} = G_2(\vec{x};\vec{\gamma}); \\ \sigma_{\phi} = G_2(\vec{x};\vec{\beta}); \\ \sigma_{\phi} = G_3(\vec{x};\vec{\delta}); \end{array}$$

- **G2**, G3, γ, δ predetermined by Monte Carlo single muon simulations and ~ negligible.
- $G1 \rightarrow ansatz$ function for the resolution on the muon momentum
- **\beta \rightarrow \beta** parameters of the p_T resolution, to be extracted from the likelihood minimization

Example of possible biases



- Dimuon mass probability before scale correction and resolution measurement
- Observed dimuon mass before scale correction
- Observed dimuon mass after scale correction
- Dimuon mass probability after momentum scale correction and resolution measurement

Strong bias in muon momentum scale evident from dimuon mass distribution